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## Structure Reports

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## Key indicators

Single-crystal X-ray study
$T=150 \mathrm{~K}$
Mean $\sigma(\mathrm{Pb}-\mathrm{Cl})=0.003 \AA$
$R$ factor $=0.019$
$w R$ factor $=0.057$
Data-to-parameter ratio $=21.4$
For details of how these key indicators were automatically derived from the article, see http://journals.iucr.org/e.
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## Lead(IV) chloride at 150 K

Lead(IV) chloride, $\mathrm{PbCl}_{4}$, is a volatile air-sensitive liquid under ambient conditions. The crystal structure consists of tetrahedral molecules residing on crystallographic twofold axes, with no intermolecular interactions within the sum of van der Waals radii. Unlike the lighter Group 14 analogues, which are based on hexagonal close packing of Cl atoms, the Cl atoms in $\mathrm{PbCl}_{4}$ form a cubic close-packed array.

## Comment

Lead(IV) chloride exists under ambient conditions as a volatile air-sensitive liquid, with a melting point of 258 K (Biltz \& Meinecke, 1923). The compound is thermally fragile and decomposes readily under ambient conditions to give lead(II) chloride and chlorine. A gas electron-diffraction study of the vapour has shown the presence of tetrahedral molecules with a $\mathrm{Pb}-\mathrm{Cl}$ bond length of 2.373 (3) $\AA$ (Haaland et al., 1992). Vibrational spectra of the liquid and solid also suggest a structure composed of essentially isolated tetrahedral molecules (Clark \& Willis, 1971; Clark \& Hunter, 1971), in contrast to the pseudo-octahedral geometries found for lead in $\mathrm{PbF}_{4}$ (Bork \& Hoppe, 1996).

(I)

The crystal structures of the Group 14 tetrachlorides, $M \mathrm{Cl}_{4}$, have been reported for $M=\mathrm{C}$ (Piermarini, 1973), Si (Zakharov et al., 1986), Ge (Merz \& Driess, 2002) and Sn (Reuter \& Pawlak, 2000). They form an isostructural series, being monoclinic, $P 2{ }_{1} / c$, with cell dimensions in the ranges $a=$ $9.07-9.86, b=5.76-6.68$ and $c=9.20-9.94 \AA$, which all increase as the group is descended from carbon to tin. All structures have values of $\beta$ in the range $103-105^{\circ}$. The structures consist of isolated tetrahedral $M \mathrm{Cl}_{4}$ molecules.

The structure of lead(IV) chloride also consists of essentially isolated tetrahedral molecules of $\mathrm{PbCl}_{4}$ (Fig. 1). However, though the crystal system is monoclinic, the lattice type is centred ( $I$-centred in the setting given here) rather than primitive, and it is clear that this compound does not follow the trend established by the lighter members of Group 14. The molecule resides on a crystallographic twofold axis at $(3 / 4, y, 0)$, where $y$ is very close to $1 / 4$. This introduces a pseudo- $C$-centring operation into the lead positions, with the

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Figure 1
The structure of $\mathrm{PbCl}_{4}$ in the crystal. The view is along the $b$ axis, and the ellipsoids enclose $50 \%$ probability surfaces.
result that, though the mean $I / \sigma$ for the data set as a whole is 27.1, that for data with $h+k \neq 2 n$ is only 5.4. There is also pseudosymmetry in the Cl positions, and this is discussed in more detail below. The $\mathrm{Pb}-\mathrm{Cl}$ bond lengths are 2.360 (2) and 2.363 (2) $\AA$, which correspond to exactly one bond-valence unit (Brown, 2002).

The shortest $\mathrm{Cl} \cdots \mathrm{Cl}$ contact observed in this structure is 3.661 (5) $\AA$, which is slightly larger than twice the van der Waals radius of chlorine ( $3.5 \AA$; Bondi, 1964). The $\mathrm{Pb} \cdots \mathrm{Cl}$ contacts are well outside the sum of the van der Waals radii of Pb and $\mathrm{Cl}(4.09 \AA)$. There are concentric shells of Cl atoms around the lead atom, the first of them consisting of 12 Cl atoms at a distance of 4.340-4.455 (2) $\AA$. The shortest $\mathrm{Pb} \cdots \mathrm{Pb}$ distance is 5.359 (1) $\AA$. As in the structures of the lighter analogues, therefore, there are no intermolecular contacts of any significance.

The structures of carbon, silicon, germanium and tin tetrachlorides can all be considered to be based on hexagonal close packing (h.c.p.) of Cl atoms (Reuter \& Pawlak, 2000). Fig. 2 is taken from the structure of $\mathrm{SnCl}_{4}$ and shows the $A B A$ layers characteristic of h.c.p. arrays. By contrast the Cl atoms in $\mathrm{PbCl}_{4}$ form a cubic close-packed (c.c.p.) array; Fig. 3 shows the characteristic $A B C$ layering sequence. Deviations from ideal symmetry may be gauged using the continuous symmetry measure (CSM) parameter described by Pinsky \& Avnir (1998). The collection of 13 atoms shown in Fig. 3 has a CSM of $0.044 \%$ relative to perfect c.c.p.; the corresponding figure for the geometry about Cl 2 is $0.037 \%$. For comparison, the CSM values for the four independent Cl atoms in $\mathrm{SnCl}_{4}$ versus perfect h.c.p. are $0.103,0.070,0.073$ and $0.081 \%$ for Cl 1 to Cl 4 , respectively. The agreement with perfect close packing is excellent in both cases; the slightly larger deviations for perfect h.c.p. are to be expected, because the layer separation in h.c.p. is not restricted by symmetry.


Figure 2
The hexagonal close packing of the Cl atoms in $\mathrm{SnCl}_{4}$. The geometry around atom Cl 1 is shown.


Figure 3
The cubic close packing of the Cl atoms in $\mathrm{PbCl}_{4}$. The geometry around atom Cl 1 is shown.

## Experimental

The sample of $\mathrm{PbCl}_{4}$ was prepared by the reaction between pyridinium hexachloroplumbate(IV) and sulfuric acid according to the method of Kauffman et al. (1983).

## Crystal data

$$
\begin{aligned}
& \mathrm{PbCl}_{4} \\
& M_{r}=349.01 \\
& \text { Monoclinic, } I 2 / a \\
& a=10.542(8) \AA \\
& b=5.359(3) \AA \\
& c=11.958(5) \AA \\
& \beta=115.83(2)^{\circ} \\
& V=608.0(6) \AA^{3} \\
& Z=4 \\
& D_{x}=3.812 \mathrm{Mg} \mathrm{~m}^{-3}
\end{aligned}
$$

Mo $K \alpha$ radiation
Cell parameters from 88 reflections
$\theta=13-21^{\circ}$
$\mu=29.35 \mathrm{~mm}^{-1}$
$T=150 \mathrm{~K}$
Cylinder, colourless
1.00 mm (length) $\times 0.20 \mathrm{~mm}$ (diameter)

## Data collection

| Stoe Stadi-4 diffractometer | $R_{\text {int }}=0.07$ |
| :--- | :--- |
| $\omega / \theta$ scans | $\theta_{\max }=25.0^{\circ}$ |
| Absorption correction: empirical | $h=-11 \rightarrow 12$ |
| $\quad$ via $\psi$ scans $[X P R E P$ (Sheldrick | $k=-6 \rightarrow 6$ |
| (2001) based on method of North | $l=-14 \rightarrow 2$ |
| et al. $(1968)]$ | 3 standard reflections |
| $T_{\min }=0.001, T_{\max }=0.002$ | every 0 reflections |
| 2322 measured reflections | frequency: 60 min |
| 536 independent reflections | intensity decay: none |

363 reflections with $I>4 \sigma(I)$

## Refinement

| Refinement on $F^{2}$ | $(\Delta \sigma)_{\max }<0.001$ |
| :--- | :--- |
| $R\left[F^{2}>2 \sigma\left(F^{2}\right)\right]=0.019$ | $\Delta \rho_{\max }=1.64 \mathrm{e}^{-3}$ |
| $w R\left(F^{2}\right)=0.057$ | $\Delta \rho_{\min }=-0.71 \mathrm{e} \AA^{-3}$ |
| $S=0.94$ | Extinction correction: Larson |
| 534 reflections | (1970), equation 22 |
| 25 parameters | Extinction coefficient: $51(3)$ |
| Weighting scheme: see below |  |

Weighting scheme: see below

## Table 1

Selected geometric parameters ( $\left(\AA,{ }^{\circ}\right.$ ).

| $\mathrm{Pb} 1-\mathrm{Cl} 1$ | $2.360(2)$ | $\mathrm{Pb} 1-\mathrm{Cl} 2$ | $2.363(2)$ |
| :--- | :--- | :--- | :--- |
|  |  |  |  |
| $\mathrm{Cl} 1-\mathrm{Pb} 1-\mathrm{Cl} 1^{\mathrm{i}}$ | $110.01(12)$ | $\mathrm{Cl} 1-\mathrm{Pb} 1-\mathrm{Cl}^{\mathrm{i}}$ | $110.35(8)$ |
| $\mathrm{Cl} 1-\mathrm{Pb} 1-\mathrm{Cl} 2$ | $107.86(8)$ | $\mathrm{Cl} 2-\mathrm{Pb} 1-\mathrm{Cl} 2^{\mathrm{i}}$ | $110.41(12)$ |

Symmetry code: (i) $\frac{3}{2}-x, y,-z$.
A stable solid-liquid equilibrium was established at 262 K in a sample of lead(IV) chloride held in a capillary mounted on a Stoe Stadi-4 diffractometer equipped with an Oxford Cryosystems lowtemperature device. A crystal was then obtained by cooling the sample at a rate of $10 \mathrm{~K} \mathrm{~h}^{-1}$. Data were collected at 150 K , a previous experiment having shown that the sample decomposed in the X-ray beam during data collection at 220 K .

The length of a crystal is difficult to control when growing a crystal of a low-melting compound in a capillary. It was almost certainly larger than the 0.8 mm collimator used. The value of $\mu$ is high, and a numerical absorption correction would normally be required. However, the cylindrical symmetry of the axially mounted sample meant that absorption anisotropy was low. A cylindrical absorption correction was attempted, and the results were satisfactory $[R(F)=$ $\left.2.3 \%, \Delta \rho_{\max }=+1.8 \mathrm{e}^{\AA^{-3}}\right]$. However, slightly better results were obtained by applying $\psi$ scans, and these are the results presented here.

The highest difference Fourier peak was $0.935 \AA$ from Pb 1 . A Tukey \& Prince (Carruthers \& Watkin, 1979) weighting scheme was
used, with $w=[$ weight $]\left[1-(\Delta F / 6 \sigma F)^{2}\right]^{2}$, where [weight] is given by a four-term Chebychev polynomial with coefficients 43.7, 56.2, 32.6 and 13.2.

Continuous symmetry measures were calculated using a locally written program based on the procedure given by Pinsky \& Avnir (1998).

The data for calculating bond valences were taken from http:// www.ccp14.ac.uk/ccp/web-mirrors/i_d_brown/bond_valence_parm/.

Data collection: DIF4 (Stoe \& Cie, 1990); cell refinement: DIF4; data reduction: REDU4 (Stoe \& Cie, 1990); program(s) used to solve structure: DIRDIF96 (Beurskens et al., 1996); program(s) used to refine structure: CRYSTALS (Watkin et al., 2001); molecular graphics: CAMERON (Watkin et al., 1996) and XP (Sheldrick, 2001); software used to prepare material for publication: CRYSTALS.

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